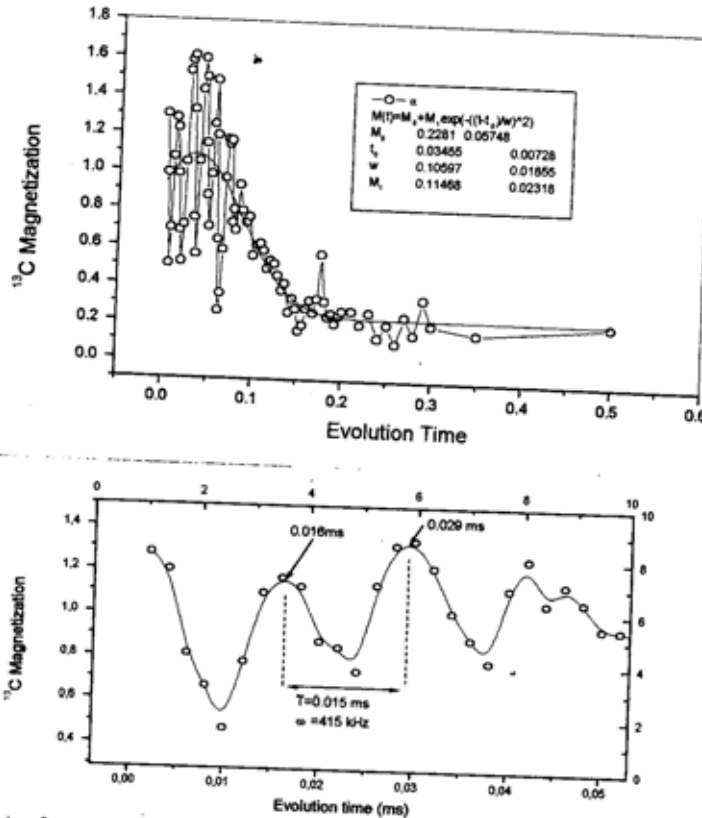


#### 4. APPENDIX 3: TIME-DOMAIN INTERFERENCES BETWEEN $^1\text{H}$ AND $^{13}\text{C}$ POLARIZATION STATES [PAS96].

We saw, in the study of the spin forward dynamics, that near the top of the echo oscillations of characteristic pulsation  $\omega_1$  are superimposed with the experimental curve.



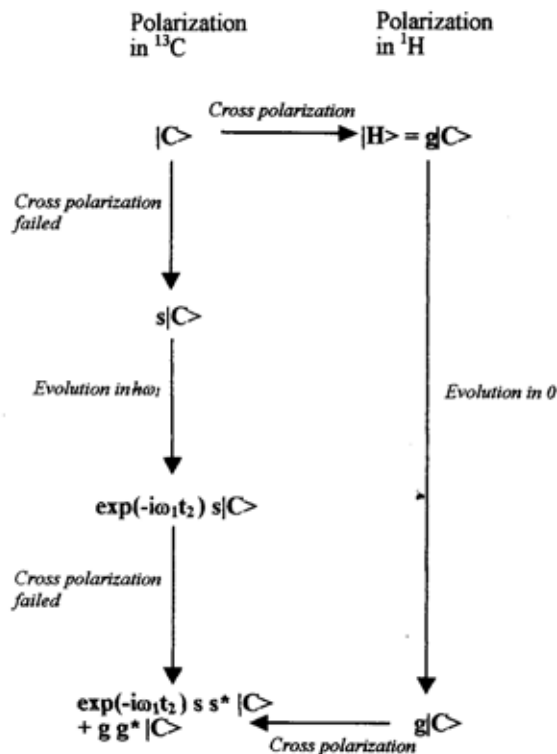
The spin forward dynamics in 5CB, and a zoom at short times. Well-defined oscillations of pulsation  $\omega_1$  appear. They can be explained as quantum interference between  $^1\text{H}$  and  $^{13}\text{C}$  polarization states.

These oscillations can be explained as follows:

During the short cross-polarizations, the polarization transfer is not perfect. (A polarization is transferred from  $^{13}\text{C}$  to  $^1\text{H}$  with a probability amplitude  $g$ , and remains in the initially polarized state with a probability amplitude  $s$  ( $|s|^2 + |g|^2 = 1$ )).

Thence the final polarization detected in  $^{13}\text{C}$  is the sum of:

- a polarization that remained all the time in  $^{13}\text{C}$ , evolving in the time  $t_2$  in the Zeeman hamiltonian  $\hbar\omega_1$  in the rotating frame, thus acquiring an additional phase  $\exp(-\frac{it_2}{\hbar}\hbar\omega_1) = \exp(-i\omega_1 t_2)$ . (The cross-polarization has never been efficient)
- a polarization that has first been transferred from  $^{13}\text{C}$  to  $^1\text{H}$ , then evolved in the time  $t_2$  in the Zeeman hamiltonian 0 in the rotating frame, then transferred from  $^1\text{H}$  to  $^{13}\text{C}$ . (Both cross-polarizations have been efficient).



The two paths of the polarization during the forward-dynamics pulse sequence  
Initially in  $^{13}\text{C}$ , it may:  
- stay in  $^{13}\text{C}$  all the time;  
- be transferred to  $^1\text{H}$  and then back to  $^{13}\text{C}$ .  
These two paths are not equivalent, since the hamiltonian in  $^{13}\text{C}$  and the one in  $^1\text{H}$  are different.  
Thus the respective wave functions acquire a relative phase and interfere with each other after rejoining.

The probability of measuring the polarization in  $^{13}\text{C}$  is, in the end:

$$P(^{13}\text{C}) = |s|^4 + |g|^4 + 2|s|^2|g|^2 \cos(\omega_1 t_2)$$

We observe very well defined *time-domain quantum interferences*.

*Remarks:* 1) The described simplistic description of the "polarization" as wave functions is of the type proposed to implement simple computations [Sci270]. It seems to work well !

2) The relative amplitude of the oscillations,  $(2|s|^2|g|^2)/(|s|^4 + |g|^4)$ , is a measure of the cross-polarization efficiency ( $|g|^2$ ).

In our experiments the efficiency is (in %):

-in 5CB: 81%

-in ferrocene: 85%